

Minicourse on fractal uncertainty principle

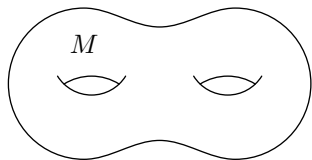
Lecture 1: overview and geometric control

Semyon Dyatlov (MIT)

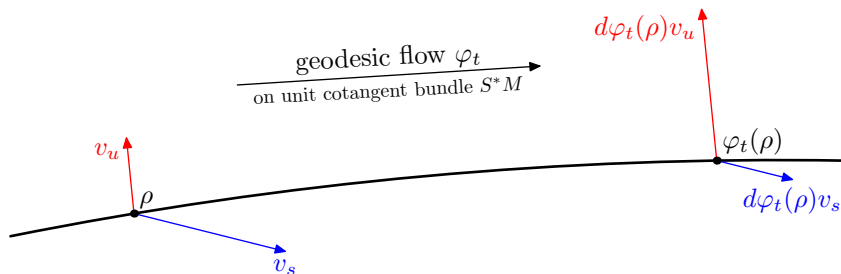
May 31, 2021

Lower bound on mass

- (M, g) compact negatively curved surface
- Geodesic flow on M : a standard model of classical chaos (perturbations diverge exponentially from the original geodesic)

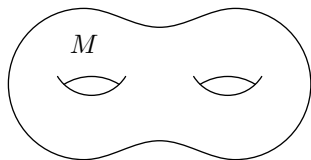


More precisely, we have the **stable/unstable** decomposition:



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$$(-\Delta_g - \lambda^2)u = 0, \quad \|u\|_{L^2} = 1$$

Theorem 1

Let $\Omega \subset M$ be a nonempty open set. Then there exists c depending on M, Ω but not on λ such that

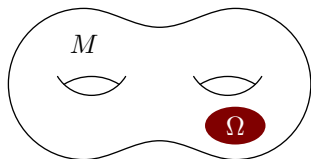
$$\|u\|_{L^2(\Omega)} \geq c > 0$$

Constant curvature: D–Jin '18, using D–Zahl '16 and Bourgain–D '18

Variable curvature: D–Jin–Nonnenmacher '19, using Bourgain–D '18

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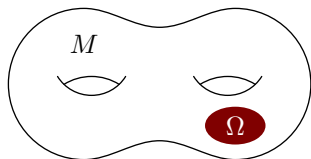
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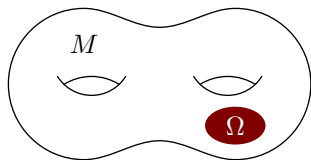
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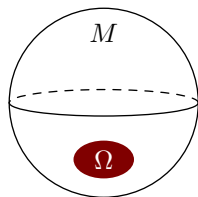
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For bounded λ this follows from unique continuation principle

The new result is in the **high frequency limit** $\lambda \rightarrow \infty$

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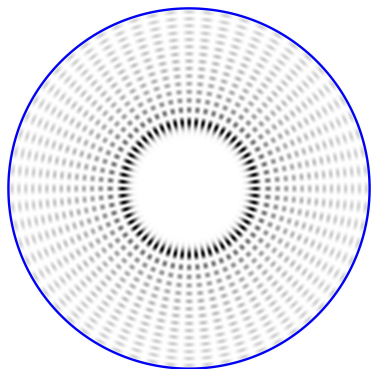
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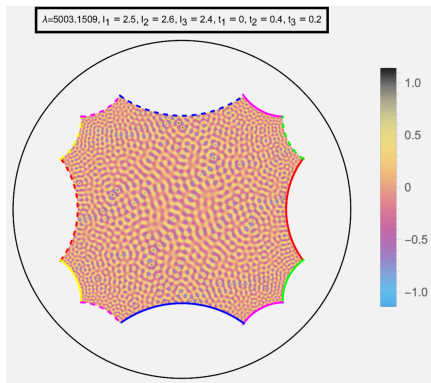
The chaotic nature of geodesic flow is important
 For example, Theorem 1 is false if M is the round sphere

An illustration

Picture on the right courtesy of Alex Strohmaier, using [Strohmaier–Uski '12](#)



Disk (Dirichlet b.c.)
Whitespace in the middle



Hyperbolic surface
No whitespace

Application to control theory [Jin '18]

Fix $T > 0$ and nonempty open $\Omega \subset M$. Then there exists $C = C(T, \Omega)$:

$$\|f\|_{L^2(M)}^2 \leq C \int_0^T \int_{\Omega} |e^{it\Delta_g} f(x)|^2 dx dt \quad \text{for all } f \in L^2(M)$$

Control by **any** nonempty open set previously known only for flat tori:
[Haraux '89](#), [Jaffard '90](#), [Burq–Zworski '04](#), [Anantharaman–Macia '14](#)...
[Datchev–Jin](#) WIP: an estimate on $C(T, \Omega)$, using [Jin–Zhang '17](#)

Application to damped wave equation [Jin '17, D–Jin–Nonnenmacher '19]

Assume that $b \in C^\infty(M)$, $b \geq 0$, $b \not\equiv 0$. Then $\exists \nu > 0$: every solution

$$v \in C^\infty([0, \infty) \times M), \quad (\partial_t^2 + b(x)\partial_t - \Delta_g)v(t, x) = 0$$

to the damped wave equation has exponentially decaying energy:

$$\int_M |\partial_t v|^2 + |\nabla_x v|^2 dx = \mathcal{O}(e^{-\nu t}) \quad \text{as } t \rightarrow \infty$$

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Microlocal analysis

Localization in **position** and **frequency** using **semiclassical quantization**

$$a(x, \xi) \in C^\infty(T^*M) \mapsto \text{Op}_h(a) = a\left(x, \frac{h}{i}\partial_x\right) : C^\infty(M) \rightarrow C^\infty(M)$$

Properties of quantization in the **semiclassical limit** $h \rightarrow 0$

- $\text{Op}_h(a) \text{Op}_h(b) = \text{Op}_h(ab) + \mathcal{O}(h)$
- $\text{Op}_h(a)^* = \text{Op}_h(\bar{a}) + \mathcal{O}(h)$
- $[\text{Op}_h(a), \text{Op}_h(b)] = -ih \text{Op}_h(\{a, b\}) + \mathcal{O}(h^2)$
- $\sup |a| < \infty \implies \|\text{Op}_h(a)\|_{L^2 \rightarrow L^2} = \mathcal{O}(1)$
- $\text{supp } b \subset \{a \neq 0\} \implies \|\text{Op}_h(b)u\| \leq C \|\text{Op}_h(a)u\| + \mathcal{O}(h^\infty)\|u\|$

$$\text{Rescale } (-\Delta_g - \lambda^2)u = 0, \quad \lambda \rightarrow \infty$$

$$\text{to obtain } (-h^2\Delta_g - 1)u = 0, \quad h = \lambda^{-1} \rightarrow 0$$

$$\text{where } -h^2\Delta_g - 1 = \text{Op}_h(p^2 - 1), \quad p(x, \xi) = |\xi|_g$$

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to obtain $(-h^2\Delta_g - 1)u = 0, \quad h = \lambda^{-1} \rightarrow 0$

where $-h^2\Delta_g - 1 = \text{Op}_h(\rho^2 - 1), \quad \rho(x, \xi) = |\xi|_g$

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Microlocal version of Theorem 1

Define the cosphere bundle $S^*M := \{(x, \xi) \in T^*M : |\xi|_g = 1\}$

Theorem 1'

Let $a \in C_c^\infty(T^*M)$ satisfy $a|_{S^*M} \not\equiv 0$. Then for $h \ll 1$ and all $u \in L^2(M)$

$$\|u\| \leq C \|\text{Op}_h(a)u\| + \frac{C \log(1/h)}{h} \|(-h^2 \Delta_g - 1)u\|$$

where the constant C depends only on M, a , but not on h, u

Remarks

- If $(-h^2 \Delta_g - 1)u = 0$ then we get $\|\text{Op}_h(a)u\| \geq c\|u\|$ for some $c > 0$
- Implies Theorem 1: $a = a(x) \implies \text{Op}_h(a)u = au$
- Sharp: $a|_{S^*M} \equiv 0, (-h^2 \Delta_g - 1)u = 0 \implies \|\text{Op}_h(a)u\| \leq Ch\|u\|$
- Cannot work for $\mathcal{O}(h/\log(1/h))$ quasimodes: Brooks '15, Eswarathasan–Nonnenmacher '17, Eswarathasan–Silberman '17

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Semiclassical measures

Take a high frequency sequence of Laplacian eigenfunctions

$$(-h_j^2 \Delta_g - 1)u_j = 0, \quad \|u_j\|_{L^2(M)} = 1, \quad h_j \rightarrow 0$$

We say u_j **converges weakly** to a measure μ on T^*M if

$$\forall a \in C_c^\infty(T^*M) : \quad \langle \text{Op}_{h_j}(a)u_j, u_j \rangle_{L^2} \rightarrow \int_{T^*M} a d\mu \quad \text{as } j \rightarrow \infty$$

Call such limits μ **semiclassical measures**

Basic properties

- μ is a probability measure, $\text{supp } \mu \subset S^*M$
- μ is invariant under the geodesic flow $\varphi_t : S^*M \rightarrow S^*M$
- Natural candidate: Liouville measure $\mu_L \sim d\text{vol}$ (equidistribution)
- Natural enemy: delta measure δ_γ on a closed geodesic (scarring)

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$$\text{Theorem 1': } a|_{S^*M} \not\equiv 0 \implies \| \text{Op}_{h_j}(a)u_j \|_{L^2} \geq c > 0$$

Theorem 1''

Let μ be a semiclassical measure on M . Then $\text{supp } \mu = S^*M$

Brief overview of history

- Quantum Ergodicity [Shnirelman '74, Zelditch '87, Colin de Verdière '85]: $\mu = \mu_L$ for density 1 sequence of eigenfunctions
- Quantum Unique Ergodicity conjecture [Rudnick–Sarnak '94]: $\mu = \mu_L$ for all eigenfunctions, that is μ_L is the only semiclassical measure. Proved in the arithmetic case [Lindenstrauss '06]

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Brief overview of history, continued

- Entropy bound [Anantharaman '08, A–Nonnenmacher '07]:
 $H_{\text{KS}}(\mu) \geq \frac{1}{2}$, in particular $\mu \neq \delta_\gamma$. Here H_{KS} denotes Kolmogorov–Sinai entropy. Note $H_{\text{KS}}(\mu_L) = 1$ and $H_{\text{KS}}(\delta_\gamma) = 0$
- Theorem 1'': between QE and QUE and 'orthogonal' to entropy bound. There exist φ_t -invariant μ with $\text{supp } \mu \neq S^*M$, $H_{\text{KS}}(\mu) > \frac{1}{2}$

Proof of Theorem 1': first steps

- In the remainder of these lectures, we will 'prove' Theorem 1'
- We focus on the case when $(-h^2\Delta_g - 1)u = 0$ and show

$$a|_{S^*M} \neq 0 \implies \|u\| \leq C_a \|\text{Op}_h(a)u\| \quad \text{for } h \ll 1$$

- We also assume that (M, g) is a **hyperbolic** surface ($K = -1$)

Partition of unity

- Take $a_1, a_2 \in C_c^\infty(T^*M \setminus 0; [0, 1])$ such that

$$a_1 + a_2 = 1 \text{ near } S^*M, \quad \text{supp } a_1 \subset \{a \neq 0\}, \quad S^*M \setminus \text{supp } a_j \neq \emptyset$$

- Define $A_j := \text{Op}_h(a_j)$. Then $A_1 + A_2 = I$ microlocally near S^*M
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Control and propagation

By ellipticity, since $\text{supp } a_1 \subset \{a \neq 0\}$,

$$\|A_1 u\| \leq C \| \text{Op}_h(a) u \| + \mathcal{O}(h^\infty) \|u\|$$

Conjugate by the half-wave propagator: for $A : L^2(M) \rightarrow L^2(M)$ and $t \in \mathbb{R}$

$$A(t) := U(-t) A U(t), \quad U(t) := \exp(-it\sqrt{-\Delta_g})$$

Since $(-h^2 \Delta_g - 1)u = 0$, we have $U(t)u = e^{-it/h}u$ and thus

$$\|A_1(t)u\| = \|A_1 u\| \leq C \| \text{Op}_h(a) u \| + \mathcal{O}(h^\infty) \|u\|$$

Egorov's Theorem: if $a \in C_c^\infty(T^*M \setminus 0)$ and t is bounded then

$$U(-t) \text{Op}_h(a) U(t) = \text{Op}_h(a \circ \varphi_t) + \mathcal{O}(h)$$

where $\varphi_t = e^{tH|_{\xi|_g}} : T^*M \setminus 0 \rightarrow T^*M \setminus 0$ is the geodesic flow

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where $\varphi_t = e^{tH|_{S^1g}} : T^*M \setminus 0 \rightarrow T^*M \setminus 0$ is the geodesic flow

Control and propagation

By ellipticity, since $\text{supp } a_1 \subset \{a \neq 0\}$,

$$\|A_1 u\| \leq C \|\text{Op}_h(a)u\| + \mathcal{O}(h^\infty)\|u\|$$

Conjugate by the half-wave propagator: for $A : L^2(M) \rightarrow L^2(M)$ and $t \in \mathbb{R}$

$$A(t) := U(-t)AU(t), \quad U(t) := \exp(-it\sqrt{-\Delta_g})$$

Since $(-h^2\Delta_g - 1)u = 0$, we have $U(t)u = e^{-it/h}u$ and thus

$$\|A_1(t)u\| = \|A_1 u\| \leq C \|\text{Op}_h(a)u\| + \mathcal{O}(h^\infty)\|u\|$$

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Sketch of the proof of Egorov's Theorem

- $a \in C_c^\infty(T^*M \setminus 0)$, $U(t) := \exp(-\frac{it}{h}P)$, $P := \sqrt{-h^2\Delta_g}$
- $\varphi_t := e^{tH_p}$ is the Hamiltonian flow of $p(x, \xi) = |\xi|_g$
- Microlocally on $T^*M \setminus 0$, $P = \text{Op}_h(p) + \mathcal{O}(h)$
- Define $A_t := \text{Op}_h(a \circ \varphi_t)$, then, since $\partial_t(a \circ \varphi_t) = \{p, a \circ \varphi_t\}$,

$$[P, A_t] = -ih \text{Op}_h(\{p, a \circ \varphi_t\}) + \mathcal{O}(h^2) = -ih\partial_t A_t + \mathcal{O}(h^2)$$

- Now $A_0 = \text{Op}_h(a)$ and

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So $U(t)A_tU(-t) = \text{Op}_h(a) + \mathcal{O}(h)$, giving Egorov's Theorem:

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Words

- Recall: $A_1 = \text{Op}_h(a_1)$, $A_2 = \text{Op}_h(a_2)$, $A(t) := U(-t)AU(t)$
- Words: $\mathcal{W}(N) := \{\mathbf{w} = w_0 \dots w_{N-1} \mid w_0, \dots, w_{N-1} \in \{1, 2\}\}$
- For $\mathbf{w} \in \mathcal{W}(N)$, define

$$A_{\mathbf{w}} := A_{w_{N-1}}(N-1) \cdots A_{w_1}(1)A_{w_0}(0), \quad a_{\mathbf{w}} := \prod_{j=0}^{N-1} (a_{w_j} \circ \varphi_j)$$

$$\text{Egorov's Theorem} \implies A_{\mathbf{w}} = \text{Op}_h(a_{\mathbf{w}}) + \mathcal{O}_N(h)$$

- We assumed $A_1 u + A_2 u = u$, so $A_1(j)u + A_2(j)u = u$ for all j . Then

$$u = \sum_{\mathbf{w} \in \mathcal{W}(N)} A_{\mathbf{w}} u$$

- Our proof will work by splitting this sum into 2 parts:
 controlled (words with enough A_1 in them) and uncontrolled, but small

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Words for the cat map (for illustration only)

$$a_{\mathbf{w}} = \prod_{j=0}^{N-1} (a_{w_j} \circ \varphi_j), \quad \mathbf{w} = w_0 \dots w_{N-1} \in \mathcal{W}(N)$$

Imagine that a_1, a_2 were indicator functions: $a_\ell = \mathbf{1}_{V_\ell}$, $S^*M = V_1 \sqcup V_2$. Then $a_{\mathbf{w}}$ is the indicator function of the set $V_{\mathbf{w}} := \bigcap_{j=0}^{N-1} \varphi_{-j}(V_{w_j})$ and $S^*M = \bigsqcup_{\mathbf{w} \in \mathcal{W}(N)} V_{\mathbf{w}}$. What do $V_{\mathbf{w}}$ look like?

Replace φ_j by the Arnold cat map

$$\varphi : \mathbb{T}^2 \rightarrow \mathbb{T}^2, \quad \mathbb{T}^2 = \mathbb{R}^2 / \mathbb{Z}^2,$$

$$\varphi(x) = \begin{pmatrix} 2 & 1 \\ 1 & 1 \end{pmatrix} x \pmod{\mathbb{Z}^2}$$

Words for the cat map (for illustration only)

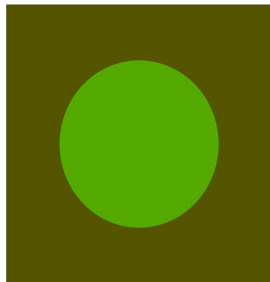
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$$N = 1$$

Replace φ_j by the **Arnold cat map**

$$\begin{aligned} \varphi : \mathbb{T}^2 &\rightarrow \mathbb{T}^2, & \mathbb{T}^2 &= \mathbb{R}^2 / \mathbb{Z}^2, \\ \varphi(x) &= \begin{pmatrix} 2 & 1 \\ 1 & 1 \end{pmatrix} x \pmod{\mathbb{Z}^2} \end{aligned}$$



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$$N = 2$$

Replace φ_j by the Arnold cat map

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Words for the cat map (for illustration only)

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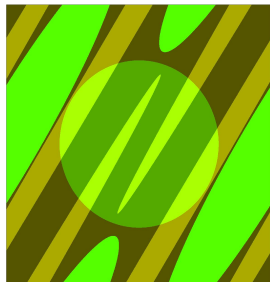
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$$N = 3$$

Replace φ_j by the Arnold cat map

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Words for the cat map (for illustration only)

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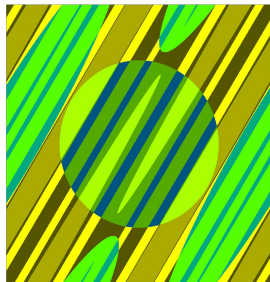
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$N = 4$

Replace φ_j by the **Arnold cat map**

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Words for the cat map (for illustration only)

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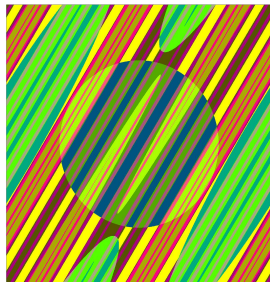
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$N = 5$

Replace φ_j by the **Arnold cat map**

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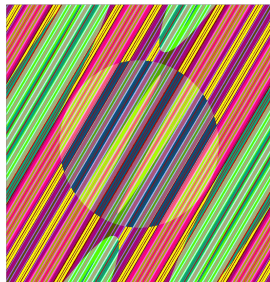
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 $N = 6$

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We see structure due to the hyperbolicity of φ
Schwarz '21: Theorem 1' for quantum cat maps

Geometric control condition (GCC)

In the rest of today's lecture we show Theorem 1':

$$(-h^2\Delta_g - 1)u = 0 \implies \|u\| \leq C \|\text{Op}_h(a)u\|$$

in the simple case when $\{a \neq 0\}$ satisfies a **geometric control condition**:

$$\exists N > 0 : S^*M \subset \bigcup_{j=0}^{N-1} \varphi_{-j}(\{a \neq 0\})$$

(This proof actually works on any compact Riemannian manifold.)

Our partition $a_1 + a_2 = 1$ has $S^*M \setminus \text{supp } a_2 \subset \text{supp } a_1 \subset \{a \neq 0\}$.

But we can choose the partition so that these sets are close to each other, so $S^*M \setminus \text{supp } a_2$ satisfies GCC. Then, taking $2 \dots 2 \in \mathcal{W}(N)$,

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Controlled/small partition under GCC

We decompose (using that $u = A_1 u + A_2 u$)

$$u = \sum_{w \in \mathcal{W}(N)} A_w u = A_x u + A_y u$$

- $A_x := A_{2\dots 2} = A_2(N-1) \cdots A_2(1)A_2(0)$ is uncontrolled
- $A_y u = \sum_{j=0}^{N-1} A_2(N-1) \cdots A_2(j+1)A_1(j)u$ where j -th term corresponds to the words w such that $w_j = 1, w_{j+1} = \cdots = w_{N-1} = 2$
- $A_x = \text{Op}_h(a_{2\dots 2}) + \mathcal{O}(h) = \mathcal{O}(h)$ since $a_{2\dots 2} = 0$ by the GCC
- $A_y u$ is estimated via $\text{Op}_h(a)u$:

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Thank you for your attention!