

Pollicott–Ruelle resonances for hyperbolic manifolds

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joint work with Frédéric Faure (Grenoble) and Colin Guillarmou (ENS)

June 25, 2014

- **Anosov flows**: hyperbolic/chaotic smooth dynamical systems
- **Pollicott–Ruelle resonances**: complex characteristic frequencies of decay of correlations
- **Our case**: geodesic flow on a hyperbolic Riemannian manifold

Resonances of the geodesic flow \sim eigenvalues of the Laplacian

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Anosov flows

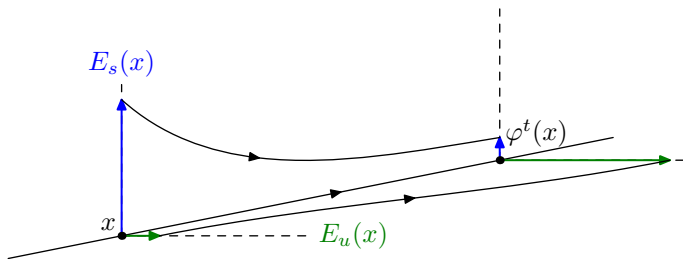
A flow $\varphi^t = e^{tX} : \mathcal{M} \rightarrow \mathcal{M}$ is **Anosov** if $(\theta > 0)$

$$T_x\mathcal{M} = \mathbb{R}X(x) \oplus E_s(x) \oplus E_u(x)$$

$$|d\varphi^t(x) \cdot v| \leq Ce^{-\theta|t|}, \quad t \geq 0, \quad v \in E_s(x)$$

$$|d\varphi^t(x) \cdot v| \leq Ce^{-\theta|t|}, \quad t \leq 0, \quad v \in E_u(x)$$

Fundamental example: $\mathcal{M} = SM$, $\varphi^t : \mathcal{M} \rightarrow \mathcal{M}$ geodesic flow,
 (M, g) has sectional curvature < 0



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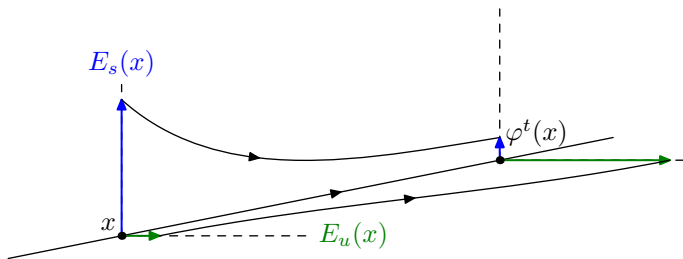
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Correlations and resonances

Assume that φ^t is **contact** \implies exists a smooth invariant measure μ

$$f, g \in L^2(\mathcal{M}) \implies \rho_{f,g}(t) = \int_{\mathcal{M}} [f \circ \varphi^{-t}] \bar{g} d\mu$$

Dolgopyat '98, Liverani '04, Tsujii '10, Nonnenmacher–Zworski '13:

$$f, g \in C^\infty \implies \rho_{f,g}(t) = \int_{\mathcal{M}} f d\mu \int_{\mathcal{M}} \bar{g} d\mu + \mathcal{O}(e^{-\varepsilon t})$$

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$$\hat{\rho}_{f,g}(\lambda) = \int_0^\infty e^{-\lambda t} \rho_{f,g}(t) dt, \quad \operatorname{Re} \lambda > 0, \quad f, g \in C^\infty$$

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$$f \circ \varphi^{-t} = e^{-tX} f, \quad X : C^\infty(\mathcal{M}) \rightarrow C^\infty(\mathcal{M}), \quad \varphi^t = e^{tX}$$

$$\hat{\rho}_{f,g}(\lambda) = \int_0^\infty \langle e^{-t(X+\lambda)} f, g \rangle_{L^2} dt = \langle (X + \lambda)^{-1} f, g \rangle_{L^2}$$

Resonances = eigenvalues of $-X$ on **anisotropic Sobolev space**

$$\mathcal{H}^r, \quad r \gg -\operatorname{Re} \lambda; \quad H^r(\mathcal{M}) \subset \mathcal{H}^r \subset H^{-r}(\mathcal{M})$$

Useful criterion: $\lambda \in \mathbb{C}$ a resonance \iff the space

$$\operatorname{Res}_X(\lambda) = \{u \in \mathcal{D}'(\mathcal{M}) \mid (X + \lambda)u = 0, \operatorname{WF}(u) \subset E_u^*\}$$

of **resonant states** is nontrivial. Here $E_u^*(x) = (\mathbb{R}X(x) \oplus E_u(x))^\perp \subset T_x^* \mathcal{M}$.

The “outgoing” condition $\operatorname{WF}(u) \subset E_u^*$ is **invariant under differential operators**

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Hyperbolic manifolds

Assumptions

- $\mathcal{M} = SM = \{(x, \xi) \mid x \in M, \xi \in T_x M, |\xi|_g = 1\}$; φ^t geodesic flow
- (M, g) has sectional curvature -1 and dimension $n + 1$
- $M = \Gamma \backslash \mathbb{H}^{n+1}$, $\Gamma \subset \text{PSO}(1, n + 1)$
- In the Minkowski hyperboloid model $\mathbb{H}^{n+1} \subset \mathbb{R}^{1, n+1}$,

$$X(x, \xi) = \xi \partial_x + x \partial_\xi$$

$$\varphi^t(x, \xi) = (x \cosh t + \xi \sinh t, x \sinh t + \xi \cosh t)$$

$$E_u(x, \xi) = \{\eta \partial_x + \eta \partial_\xi \mid \eta \in T_x \mathbb{H}^{n+1}, \eta \perp \xi\}$$

$$E_s(x, \xi) = \{\eta \partial_x - \eta \partial_\xi \mid \eta \in T_x \mathbb{H}^{n+1}, \eta \perp \xi\}$$

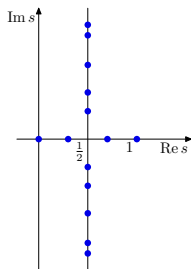
Gelfand–Fomin '55, Guillemin '77, Moore '87, Ratner '87, Zelditch '87, Flaminio–Forni '03, Leboeuf '04, Faure–Tsuji '13...

Results

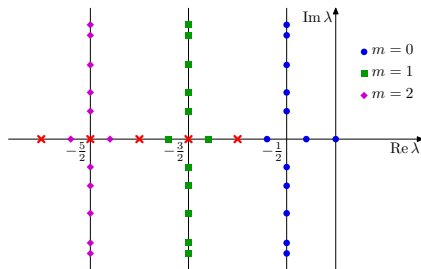
Theorem, $\dim M = 2$

Except at $-1 - \frac{1}{2}\mathbb{N}_0$, λ is a P-R resonance iff

$$-(\lambda + m)(\lambda + m + 1) \in \text{Spec}(-\Delta) \quad \text{for some } m \in \mathbb{N}_0.$$



$$s(1 - s) \in \text{Spec}(-\Delta)$$



$$\lambda \in \text{Spec}_{\text{PR}}(X)$$

Theorem [DFG '14], $\dim M = n + 1 \geq 2$

Except at $-\frac{n}{2} - \frac{1}{2}\mathbb{N}_0$ (+ special case: $\lambda \in -2\mathbb{N}$), λ is a resonance iff

$$-\left(\lambda + m + \frac{n}{2}\right)^2 + \frac{n^2}{4} + m - 2\ell \in \text{Spec}^{m-2\ell}(-\Delta), \quad m \in \mathbb{N}_0, \ell \in [0, m/2],$$

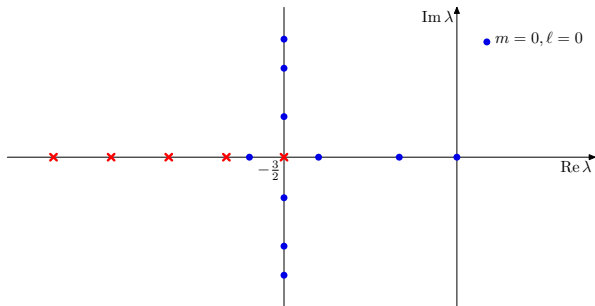
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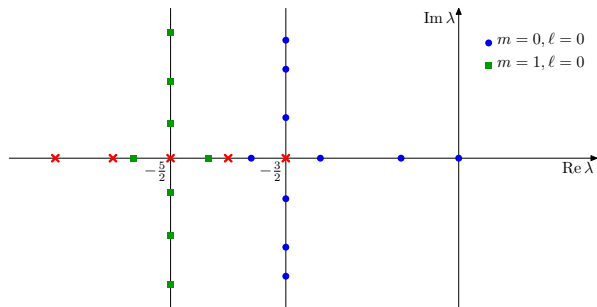


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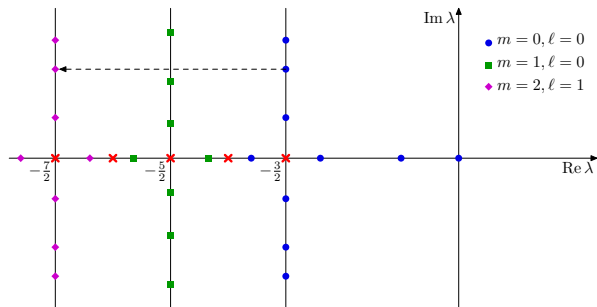


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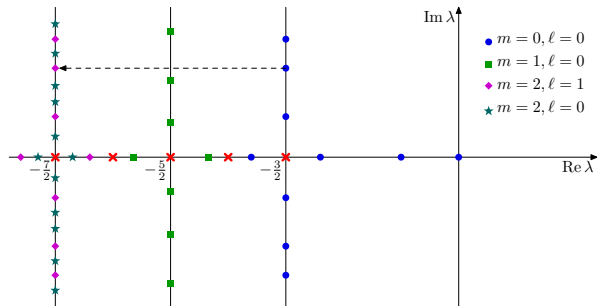


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Resonant states and multiplicities

Resonant states of $X \sim$ eigenstates of Δ

$$\text{Res}_X(\lambda) = \{u \in \mathcal{D}'(SM) \mid (X + \lambda)u = 0, \text{WF}(u) \subset E_u^*\}$$

$$\text{Res}_X^*(\lambda) = \{u^* \in \mathcal{D}'(SM) \mid (X^* + \bar{\lambda})u^* = 0, \text{WF}(u^*) \subset E_s^*\}$$

$$\text{Res}_X^\#(\lambda) = \{u \in \mathcal{D}'(SM) \mid \exists k : (X + \lambda)^k u = 0, \text{WF}(u) \subset E_u^*\}$$

Theorem [DFG '14]: resonances are semisimple

For $\lambda \notin -\frac{n}{2} - \frac{1}{2}\mathbb{N}_0$, $\text{Res}_X(\lambda) = \text{Res}_X^\#(\lambda)$

Proof relies on pairing formula:

$$\langle u, u^* \rangle_{L^2(SM)} = \mathcal{F}(\lambda) \langle f, f^* \rangle_{L^2(M)}, \quad u \in \text{Res}_X(\lambda), \quad u^* \in \text{Res}_X^*(\lambda)$$

where f, f^* are the eigenstates of $-\Delta$ on tensors which correspond to u, u^*
 Anantharaman–Zelditch '07, Hansen–Hilgert–Schröder '11

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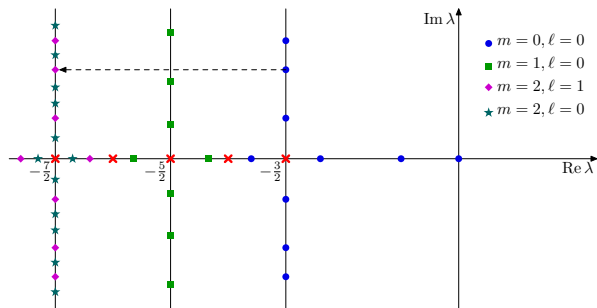
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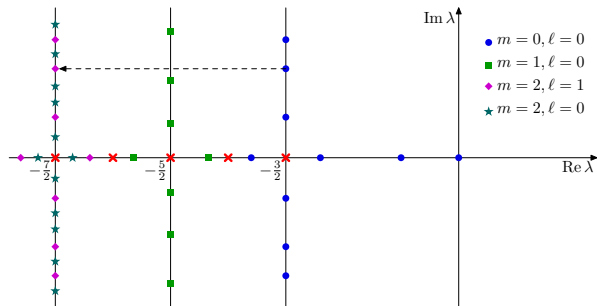
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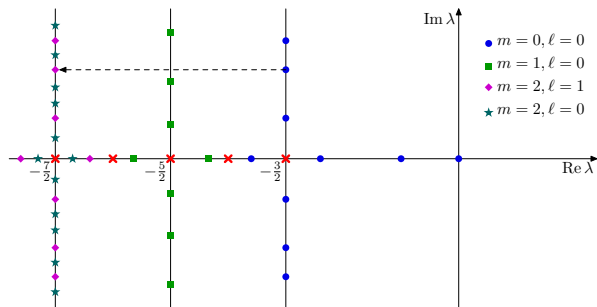
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- Rate of decay of correlations depends on the small eigenvalues of Δ
- Possible applications to quantum ergodicity of resonant states

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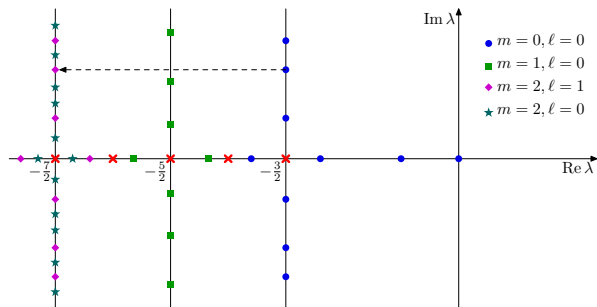
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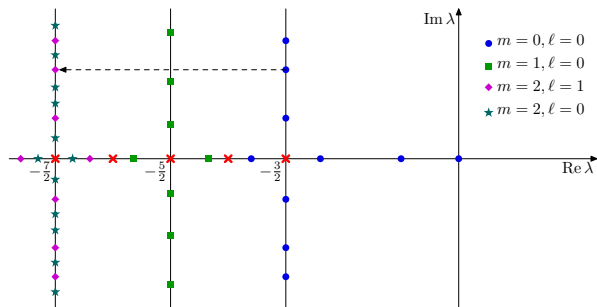
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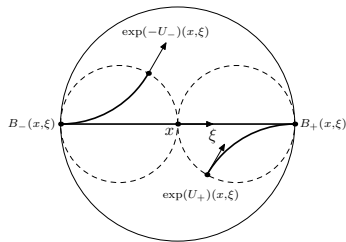
Horocyclic operations (dimension 2)

$$M = \Gamma \backslash \mathbb{H}^2, \quad SM = \Gamma \backslash S\mathbb{H}^2, \quad \varphi^t = e^{tX}$$

U_+, U_- horocyclic vector fields on SM :

$$[X, U_{\pm}] = \pm U_{\pm}, \quad [U_+, U_-] = 2X$$

$$E_u = \mathbb{R}U_-, \quad E_s = \mathbb{R}U_+$$



λ is a resonance $\iff \exists u \in \mathcal{D}'(SM) \setminus 0 : (X + \lambda)u = 0, \quad \text{WF}(u) \subset E_u^*$

$$m \geq 0 \implies (X + \lambda + m)U_-^m u = 0, \quad \text{WF}(U_-^m u) \subset E_u^*$$

No resonances for $\text{Re } \lambda > 0 \implies U_-^m u = 0$ for $m > -\text{Re } \lambda$

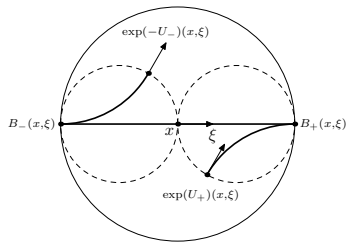
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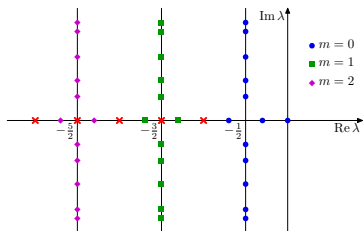
$$(X + \lambda)u = 0, \quad \text{WF}(u) \subset E_u^*, \quad U_-^{m+1}u = 0 \quad \text{for } m \gg 1$$

$$u, U_-u, U_-^2u, \dots, v := U_-^m u \neq 0, \quad U_-^{m+1}u = 0$$

$$(X + \lambda + m)v = 0, \quad U_-v = 0; \quad \text{WF}(v) \subset E_u^* \text{ automatically}$$

Other direction: given v , put $(\lambda \notin -1 - \frac{1}{2}\mathbb{N})$

$$u := \frac{(2\lambda + m)!}{m!(2\lambda + 2m)!} U_+^m v; \quad (X + \lambda)u = 0, \quad U_-^m u = v$$



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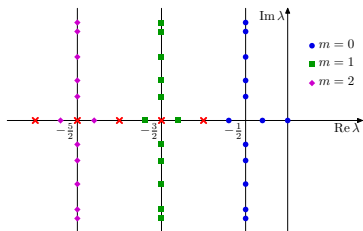
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First band (dimension 2)

$$(X + \lambda)v = 0, \quad U_- v = 0$$

$$f(x) := \int_{S_x M} v(x, \xi) d\xi$$

$$(-\Delta + \lambda(\lambda + 1))f = 0$$

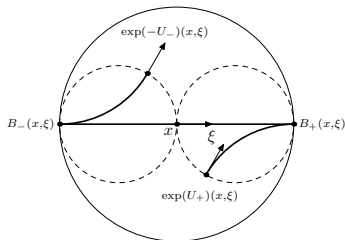
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$$U_- B_- = X B_- = 0, \quad U_- \Phi_- = 0, \quad X \Phi_- = -\Phi_-$$

$$v(x, \xi) = \Phi_-(x, \xi)^\lambda w(B_-(x, \xi)), \quad w \in \mathcal{D}'(S^1)$$

$$w(\gamma \cdot \nu) = |d\gamma(\nu)|^\lambda w(\nu), \quad \nu \in S^1, \quad \gamma \in \Gamma$$

$$f(x) = \int_{S^1} P(x, \nu)^{1+\lambda} w(\nu) d\nu, \quad P(x, \nu) = \frac{1 - |x|^2}{|x - \nu|^2}, \quad x \in \mathbb{B}^2$$



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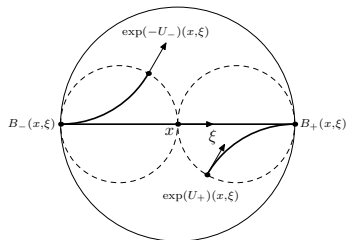
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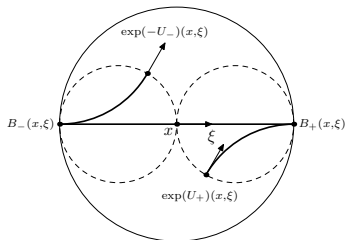
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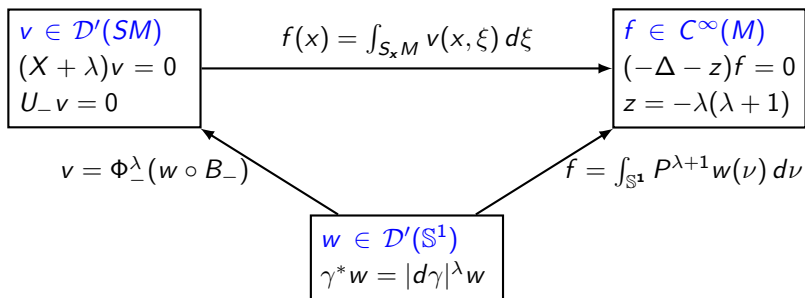
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$$w(\gamma \cdot \nu) = |d\gamma(\nu)|^\lambda w(\nu), \quad \nu \in \mathbb{S}^1, \quad \gamma \in \Gamma$$

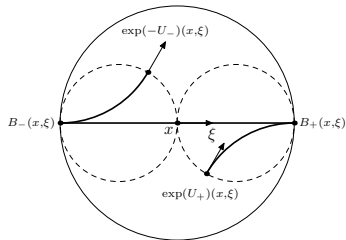
$$f(x) = \int_{\mathbb{S}^1} P(x, \nu)^{1+\lambda} w(\nu) d\nu, \quad P(x, \nu) = \frac{1 - |x|^2}{|x - \nu|^2}, \quad x \in \mathbb{B}^2$$



First band (dimension 2)

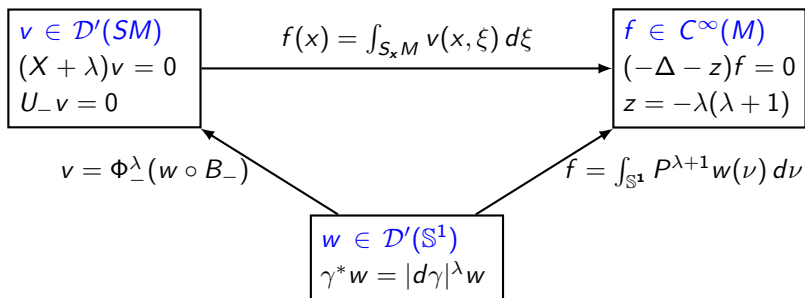


Poisson operator $w \mapsto f$ bijective ($\lambda \notin -1 - \mathbb{N}$):
 Helgason '70, '74, Minemura '75,
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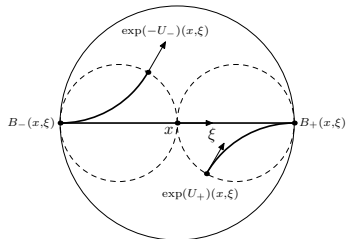


This finishes the proof in dimension 2

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This finishes the proof in dimension 2

Higher dimensions

Problem: U_{\pm} are no longer well-defined vector fields on SM

Solution: in dimension 2, U_- spanned E_u

$$E_u(x, \xi) = \theta_-(\mathcal{E}(x, \xi)), \quad \theta_- : \eta \mapsto (\eta, \eta)$$

$$\mathcal{E}(x, \xi) = \{\eta \in T_x M \mid \eta \perp \xi\} \quad \text{vector bundle over } SM$$

$$U_- : \mathcal{D}'(SM) \rightarrow \mathcal{D}'(SM; \mathcal{E}^*), \quad U_- u(x, \xi) \cdot \eta = du(x, \xi) \cdot \theta_-(\eta)$$

$$U_-^m : \mathcal{D}'(SM) \rightarrow \mathcal{D}'(SM; \otimes_S^m \mathcal{E}^*), \quad U_-^m X = (X + m)U_-^m$$

Start with $u \in \mathcal{D}'(SM)$, $(X + \lambda)u = 0$, $\text{WF}(u) \subset E_u^*$

$$u, U_- u, U_-^2 u, \dots, v := U_-^m u \neq 0, U_-^{m+1} u = 0$$

$$v \in \mathcal{D}'(SM; \otimes_S^m \mathcal{E}^*), \quad (X + \lambda + m)v = 0, \quad U_- v = 0$$

Analogy with **polynomials on \mathbb{R}^n** : $U_- \sim d$, $U_+ \sim x$

$v \mapsto u$: "Taylor formula" but **commutation relations more involved**

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Poisson on symmetric trace-free tensors ($v \in \widehat{\otimes}_S^m \Leftrightarrow g^{\alpha\beta} v_{\alpha\beta J} = 0$):

$$\mathcal{P}_\lambda : \{w \in \mathcal{D}'(\mathbb{S}^n; \widehat{\otimes}_S^m(T^*\mathbb{S}^n)) \mid \gamma^* w = |d\gamma|^{\lambda+m} w\} \rightarrow \\ \{v \in C^\infty(M; \widehat{\otimes}_S^m(T^*M)) \mid \nabla^* v = 0, (-\Delta + \lambda(n+\lambda) + m)v = 0\}$$

is an **isomorphism** for $\lambda \notin -\frac{n}{2} - \frac{1}{2}\mathbb{N}_0$, $n > 1$

Injectivity

Weak expansion ($x = r\nu \in \mathbb{B}^{n+1}$, $\rho = 1 - r \rightarrow 0$, $\nu \in \mathbb{S}^n$)

$$\mathcal{P}_\lambda w(x) \sim c_\lambda w(\nu) \rho^{-\lambda} + \sum_{k \in \mathbb{N}} f_k(\nu) \rho^{-\lambda+k} + \sum_{k \in \mathbb{N}_0} f'_k(\nu) \rho^{n+\lambda+k}$$

Surjectivity

- \exists weak expansion for each trace-free divergence-free eigenstate v
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Pairing formula and semisimplicity ($m = 0$)

$$(X + \lambda)v = 0, \quad \mathcal{U}_- v = 0, \quad f(y) = \int_{S_x M} v(y, \eta_-) d\eta_-$$

$$(X^* + \bar{\lambda})v^* = 0, \quad \mathcal{U}_+ v^* = 0, \quad f^*(y) = \int_{S_x M} v^*(y, \eta_+) d\eta_+$$

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Change of variables $\Psi : (x, \xi, \eta) \mapsto (y, \eta_-, \eta_+)$

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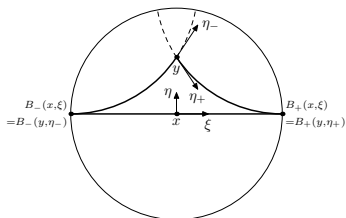
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Thank you for your attention!